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# Lower Order Eigenvalues of the Schrödinger Operator\* Research Article

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**Abstract.** Making use of the method introduced by Brands in [2], we consider lower order eigenvalues of the Schrödinger operator in Euclidean domains. We extend an estimate on eigenvalues obtained by Ashbaugh and Benguria in [1].

Keywords. Membrane eigenvalue; Schrödinger operator; Rayleigh-Ritz inequality

MSC. 35P15; 58C40

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### 1. Introduction

Let  $\Omega$  be a bounded domain in an n-dimensional Euclidean space  $\mathbb{R}^n$  with smooth boundary  $\partial\Omega$ . The eigenvalue problem

$$\begin{cases}
-\Delta u = \lambda u, & \text{in } \Omega, \\
u = 0, & \text{on } \partial\Omega
\end{cases}$$
(1.1)

is called the *fixed membrane problem*. Let  $0 < \lambda_1 < \lambda_2 \le \lambda_3 \le \cdots \to +\infty$  denote the successive eigenvalues for (1.1), where each eigenvalue is repeated according to its multiplicity. In the case of n=2, Payne-Pólya-Weinberger [5] proved

$$\frac{\lambda_2 + \lambda_3}{\lambda_1} \le 6. \tag{1.2}$$

Subsequently, in 1964, Brands [2] sharpen (1.2) to

$$\frac{\lambda_2 + \lambda_3}{\lambda_1} \le 5 + \frac{\lambda_1}{\lambda_2}.\tag{1.3}$$

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In 1993, for general dimensions  $n \ge 2$ , Ashbaugh and Benguria [1] proved (see the inequality (6.10) in [1])

$$\frac{\lambda_2 + \lambda_3 + \dots + \lambda_{n+1}}{\lambda_1} \le n + 3 + \frac{\lambda_1}{\lambda_2}. \tag{1.4}$$

Recently, the inequality (1.4) has been extended to some Riemannian manifolds, see [6, 3, 4] and the references therein.

In this note, we consider eigenvalue problem of the following Schrödinger operator

$$\begin{cases} (-\Delta + V)u = \lambda u, & \text{in } \Omega, \\ u = 0, & \text{on } \partial\Omega, \end{cases}$$
 (1.5)

where V is a continuous bounded function on  $\overline{\Omega}$ . Using the method of Brands [2], we study the eigenvalue problem (1.5) for general dimensions  $n \ge 2$  and extend the inequality (1.4) as follows:

**Theorem.** Let  $\lambda_i$  be the i-th eigenvalue of the eigenvalue problem (1.5). Then

$$\frac{\lambda_2 + \lambda_3 + \dots + \lambda_{n+1}}{\lambda_1} \le n + \frac{(M+1)(3\xi + 4M + 1)}{\xi + M},\tag{1.6}$$

where  $M = \sup_{\overline{\Omega}} |V|/\lambda_1$  and  $\xi = \lambda_2/\lambda_1$ .

**Remark.** If V = 0 in (1.6), from (1.6), it is easy to see that

$$\frac{\lambda_2 + \lambda_3 + \dots + \lambda_{n+1}}{\lambda_1} \le n + 3 + \frac{\lambda_1}{\lambda_2},$$

(1.4) follows. Hence, (1.6) extends the inequality (1.4).

### 2. Proof of Theorem

Let  $u_i$  be the orthonormal eigenvalue function with respect to  $L^2$  inner product corresponding to  $\lambda_i$ , that is,

$$\int_{\Omega} u_i u_j = \delta_{ij}, \quad \text{for any } i, j.$$

We choose rectangular coordinates  $\tilde{x}^1, \tilde{x}^2, \dots, \tilde{x}^n$  of the Euclidean space  $\mathbb{R}^n$  by taking as origin the center of gravity of  $\Omega$  with mass-distribution  $u_1^2$  such that

$$\int_{\Omega} \tilde{x}^{i} u_{1}^{2} = 0, \quad \text{for } i = 1, 2, \dots, n.$$
(2.1)

Defining an  $n \times n$ -matrix B as follows:

$$B := (b_{i,i})$$

where  $b_{ij} = \int_{\Omega} \tilde{x}^i u_1 u_{j+1}$ . Using the orthogonalization of Gram and Schmidt, we know that there exist an upper triangle matrix  $R = (R_{ij})$  and an orthogonal matrix  $Q = (q_{ij})$  such that R = QB, that is,

$$R_{ij} = \sum_{k=1}^{n} q_{ik} b_{kj} = \int_{k} \sum_{k=1}^{n} q_{ik} \widetilde{x}^{k} u_{1} u_{j} = 0, \quad 2 \le j \le i \le n.$$
(2.2)

Setting  $x^i = \sum_{j=1}^n q_{ij} \tilde{x}^j$ . From (2.1) and (2.2), we arrive at

$$\int_{\Omega} x_i u_1 u_j = 0, \quad \text{for } 1 \le j \le i \le n.$$

$$(2.3)$$

Let  $\varphi_i = x_i u_1$ . Then  $\varphi_i = 0$  on  $\partial \Omega$  and

$$\int_{\Omega} \varphi_i u_j = 0, \quad \text{for } 1 \le j \le i \le n.$$

One gets from Rayleigh-Ritz inequality that

$$\lambda_{i+1} \le \frac{\int_{\Omega} \varphi_i(-\Delta + V)\varphi_i}{\int_{\Omega} \varphi_i^2} \,. \tag{2.4}$$

Note that

$$(-\Delta + V)\varphi_i = \lambda_1 x_i u_1 - 2u_{1,x_i},$$

where  $u_{1,x_i} = \partial u_1/\partial x_i$ . It follows that

$$\int_{\Omega} \varphi_{i}(-\Delta + V)\varphi_{i} = \int_{\Omega} \varphi_{i}(\lambda_{1}x_{i}u_{1} - 2u_{1,x_{i}})$$

$$= \lambda_{1} \int_{\Omega} \varphi_{i}^{2} - 2 \int_{\Omega} x_{i}u_{1}u_{1,x_{i}}$$

$$= \lambda_{1} \int_{\Omega} \varphi_{i}^{2} - \int_{\Omega} x_{i}(u_{1}^{2})_{,x_{i}}$$

$$= \lambda_{1} \int_{\Omega} \varphi_{i}^{2} + \int_{\Omega} u_{1}^{2}$$

$$= \lambda_{1} \int_{\Omega} \varphi_{i}^{2} + 1.$$
(2.5)

(2.5) combining with (2.4) yields

$$\lambda_{i+1} \le \lambda_1 + \left( \int_{\Omega} (x_i u_1)^2 \right)^{-1}. \tag{2.6}$$

By integration by parts, it holds that

$$\int_{\Omega} u_1^{\alpha+1} = -\int_{\Omega} x_i (u_1^{\alpha+1})_{,x_i} = -(\alpha+1) \int_{\Omega} (x_i u_1) (u_1^{\alpha-1} u_{1,x_i}).$$

For  $\alpha > 1/2$ , it follows from the Cauchy-Schwarz inequality that

$$\begin{split} \left(\int_{\Omega} u_{1}^{\alpha+1}\right)^{2} &= (\alpha+1)^{2} \left(\int_{\Omega} (x_{i}u_{1})(u_{1}^{\alpha-1}u_{1,x_{i}})\right)^{2} \\ &\leq (\alpha+1)^{2} \int_{\Omega} (x_{i}u_{1})^{2} \int_{\Omega} (u_{1}^{\alpha-1}u_{1,x_{i}})^{2} \\ &= \frac{(\alpha+1)^{2}}{2\alpha-1} \int_{\Omega} (x_{i}u_{1})^{2} \int_{\Omega} (u_{1}^{2\alpha-1})_{,x_{i}} u_{1,x_{i}} \\ &= \frac{-(\alpha+1)^{2}}{2\alpha-1} \int_{\Omega} (x_{i}u_{1})^{2} \int_{\Omega} u_{1}^{2\alpha-1} u_{1,x_{i}x_{i}}. \end{split}$$

Thus

$$\left(\int_{\Omega} (x_i u_1)^2\right)^{-1} \le \frac{-(\alpha+1)^2}{2\alpha-1} \frac{\int_{\Omega} u_1^{2\alpha-1} u_{1,x_i x_i}}{\left(\int_{\Omega} u_1^{\alpha+1}\right)^2}.$$
 (2.7)

Applying (2.7) to (2.6), one gets

$$\frac{\lambda_2 + \lambda_3 + \dots + \lambda_{n+1}}{\lambda_1} \le n + \frac{(\alpha + 1)^2}{2\alpha - 1} A(\alpha) \sup_{\overline{\Omega}} \left( 1 - \frac{V}{\lambda_1} \right) 
\le n + (M+1) \frac{(\alpha + 1)^2}{2\alpha - 1} A(\alpha),$$
(2.8)

where  $A(\alpha) = \int_{\Omega} u_1^{2\alpha} / (\int_{\Omega} u_1^{\alpha+1})^2$ .

In the following, we will find an upper bound of

$$\frac{(\alpha+1)^2}{2\alpha-1}A(\alpha). \tag{2.9}$$

Define

$$\phi = u_1^{\alpha} - u_1 \int_{\Omega} u_1^{\alpha+1}$$
, for  $\alpha > 1$ .

Then we have

$$\int_{\Omega} \phi u_1 = 0.$$

This means that

$$\lambda_2 \le \frac{\int_{\Omega} \phi(-\Delta + V)\phi}{\int_{\Omega} \phi^2}.$$
 (2.10)

Note that

$$\begin{split} \alpha \int_{\Omega} u_1^{\alpha-1} |\nabla u_1|^2 &= \int_{\Omega} u_1^{\alpha} (-\Delta) u_1 \\ &= \int_{\Omega} (\lambda_1 - V) u_1^{\alpha+1}, \\ (2\alpha - 1) \int_{\Omega} u_1^{2\alpha - 2} |\nabla u_1|^2 &= \int_{\Omega} u_1^{2\alpha - 1} (-\Delta) u_1 \\ &= \int_{\Omega} (\lambda_1 - V) u_1^{2\alpha}, \end{split}$$

and

$$(-\Delta+V)\phi = -\alpha(\alpha-1)u_1^{\alpha-2}|\nabla u_1|^2 + (\alpha\lambda_1 - \alpha V + V)u_1^{\alpha} - \lambda_1 u_1 \int_{\Omega} u_1^{\alpha+1}.$$

Hence, we have

$$\int_{\Omega} \phi(-\Delta + V)\phi = \frac{\alpha^{2}}{2\alpha - 1} \lambda_{1} \int_{\Omega} u_{1}^{2\alpha} - \frac{(\alpha - 1)^{2}}{2\alpha - 1} \int_{\Omega} V u_{1}^{2\alpha} - \lambda_{1} \left( \int_{\Omega} u_{1}^{\alpha + 1} \right)^{2} \\
\leq \left( \frac{\alpha^{2}}{2\alpha - 1} + \frac{(\alpha - 1)^{2}}{2\alpha - 1} M \right) \lambda_{1} \int_{\Omega} u_{1}^{2\alpha} - \lambda_{1} \left( \int_{\Omega} u_{1}^{\alpha + 1} \right)^{2}.$$
(2.11)

From (2.10) and (2.11), we arrive at

$$\frac{\lambda_2}{\lambda_1} \le \frac{\left(\frac{\alpha^2}{2\alpha - 1} + \frac{(\alpha - 1)^2}{2\alpha - 1}M\right)A(\alpha) - 1}{A(\alpha) - 1}.$$
(2.12)

Again, by using the Cauchy-Schwarz inequality, one gets

$$\left(\int_{\Omega}u_1^{\alpha+1}\right)^2 = \left(\int_{\Omega}u_1^{\alpha}u_1\right)^2 \leq \int_{\Omega}u_1^{2\alpha}\int_{\Omega}u_1^2 = \int_{\Omega}u_1^{2\alpha}.$$

This means that  $A(\alpha) > 1$  for  $\alpha > 1$ . If  $\alpha$  is restricted to the condition

$$\xi - \left(\frac{\alpha^2}{2\alpha - 1} + \frac{(\alpha - 1)^2}{2\alpha - 1}M\right) > 0,$$

that is

$$1 < \alpha < \frac{(M+\xi) + \sqrt{(M+\xi)(\xi-1)}}{M+1}. \tag{2.13}$$

Then (2.12) is equivalent to

$$A(\alpha) \le \frac{\xi - 1}{\xi - \left(\frac{\alpha^2}{2\alpha - 1} + \frac{(\alpha - 1)^2}{2\alpha - 1}M\right)},\tag{2.14}$$

where  $\xi = \lambda_2/\lambda_1$ . Inserting (2.14) into (2.8) yields

$$\frac{\lambda_2 + \lambda_3 + \dots + \lambda_{n+1}}{\lambda_1} \le n + (M+1)(\xi - 1)f(\alpha), \tag{2.15}$$

where

$$f(\alpha) = \frac{(\alpha+1)^2}{(2\alpha-1)\xi - [\alpha^2 + (\alpha-1)^2 M]}.$$

The minimum of  $f(\alpha)$  as a function of  $\alpha$  in the range (2.13) is

$$\frac{(M+1)(3\xi + 4M + 1)}{(\xi + M)(\xi - 1)}$$

and this is attained at

$$\alpha = \frac{2\xi + 2M}{\xi + 2M + 1}.$$

Hence, (2.15) yields

$$\frac{\lambda_2 + \lambda_3 + \dots + \lambda_{n+1}}{\lambda_1} \le n + (M+1)(\xi - 1)f\left(\frac{2\xi + 2M}{\xi + 2M + 1}\right)$$
$$= n + \frac{(M+1)(3\xi + 4M + 1)}{\xi + M}.$$

This concludes the proof of theorem.

### References

- [1] M. S. Ashbaugh, R. D. Benguria, More bounds on eigenvalue ratios for Dirichlet Laplacians in *n* dimensions, *SIAM J. Math. Anal.* **24** (1993), 1622–1651.
- <sup>[2]</sup> J. J. A. M. Brands, Bounds for the ratios of the first three membrane eigenvalues, *Arch. Rational Mech. Anal.* **16** (1964), 265–268.
- D. G. Chen, Q. M. Cheng, Extrinsic estimates for eigenvalues of the Laplace operator, *J. Math. Soc. Japan* **60** (2008), 325–339.
- [4] G. Y. Huang, X. X. Li, R. W. Xu, Extrinsic estimates for the eigenvalues of Schrödinger operator, Geom. Dedicata 143 (2009), 89–107.
- <sup>[5]</sup> L. E. Payne, G. Pólya, H. F. Weinberger, On the ratio of consecutive eigenvalues, *J. Math. Phys.* **35** (1956), 289–298.
- [6] H. J. Sun, Q. M. Cheng, H. C. Yang, Lower order eigenvalues of Dirichlet Laplacian, Manuscripta Math. 125 (2008), 139–156.